

# A first principles study of hydrogen storage in lithium decorated defective phosphorene



Sandip Haldar <sup>a</sup>, Sankha Mukherjee <sup>a</sup>, Farheen Ahmed <sup>a</sup>, Chandra Veer Singh <sup>a,b,\*</sup>

<sup>a</sup> Department of Materials Science and Engineering, University of Toronto, Toronto, ON, M5S 3E4, Canada <sup>b</sup> Department of Mechanical and Industrial Engineering, University of Toronto, Toronto, ON, M5S 3E4, Canada

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### ABSTRACT

In this study, we studied defect-engineering and lithium decoration of 2D phosphorene for effective hydrogen storage using density functional theory. Contrary to graphene, it is found that the presence of point-defects is not preferable for anchoring of  $H_2$  molecules over defective phosphorene. According to previous research, strategies such as defect engineering, metal decoration, and doping enhance the hydrogen storage capacity of several 2D materials. Our DFT simulations show that point defects in phosphorene do not improve the hydrogen storage capacity compared to pristine phosphorene. However, selective lithium decoration over the defective site significantly improves the hydrogen adsorption capacity yielding a binding energy of as high as -0.48 eV/H<sub>2</sub> in Li-decorated single vacancy phosphorene. Differential charge densities and projected density of states have been computed to understand the interactions and charge transfer among the constituent atoms. Strong polarization of the  $H_2$  molecule is evidenced by the charge accumulation and depletion. The PDOS shows that the presence of Li leads to enhanced charge transfer. The maximum gravimetric density was investigated by sequentially adding  $H_2$ molecules to the Li-decorated single vacancy defective phosphorene. The Li-decorated single vacancy phosphorene is found to possess a gravimetric density of around 5.3% for hydrogen storage.

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# Introduction

Hydrogen fuel has attracted increased attention as a source of clean energy in order to tackle the adverse environmental issues of fossil fuels due to carbon emission. Fossil fuel based energy systems produce pollutants in greater amounts and are more damaging than those produced by renewable hydrogen energy systems. In fact, an engine that burns pure hydrogen produces almost no environmental pollution [1]. Therefore, hydrogen based energy systems are promising for solving environmental problems and the energy crisis [2]. Significant research effort is being invested in developing safe, cost-effective and efficient hydrogen storage media for fuel cell applications [3]. Barthelemy et al. [4] have reviewed the industrial perspective of  $H_2$  fuel cells in the light of the  $H_2$  fuel market with applications as fuel in transportation and stationary power sources.

One of the primary concerns of using  $H_2$  fuel as a practical alternative to fossil fuels is its storage and handling.

\* Corresponding author. Department of Materials Science and Engineering, University of Toronto, Toronto, ON, M5S 3E4, Canada. E-mail address: Chandraveer.Singh@utoronto.ca (C.V. Singh).

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Conventional methods to store  $H_2$  use high pressure gas cylinders (up to 800 bar) or as liquid H<sub>2</sub> in cryogenic tanks (at 21 K) [5]. However, these methods are inefficient and unsafe for automotive applications. In order to achieve efficient and safe storage, H<sub>2</sub> adsorption into substrates with a high specific surface area is a promising route [5,6]. Effective  $H_2$  storage based on physisorption mechanism requires: (a) a high gravimetric density: a target of 7.5 wt % system gravimetric density was decided as revised target the United States Department of Energy (USDOE) [7] for ultimate full fleet applications, and (b) easy recycling of hydrogen at ambient pressure and temperature [8,9]. While experimental studies are costly, computational investigations have emerged as a viable avenue to determine the performance of a potential host substrate for H<sub>2</sub> storage. In the past, density functional theory (DFT) calculations were extensively used to investigate the possibility and capacity of a material for effective H<sub>2</sub> storage. The leading parameter to identify the potential of a material for H<sub>2</sub> storage is its binding energy that should be in the range of between 0.2 and 0.6 eV/H<sub>2</sub> for reversible H<sub>2</sub> storage (i.e. adsorption and release at near ambient conditions) [10]. The maximum gravimetric density can be predicted by sequential adsorption of  $H_2$  molecules and relating it to the binding energy of the system.

Several materials have been demonstrated to be potential substrates for  $H_2$  storage, e.g. metal organic frameworks [11], covalent-organic frameworks [12], carbon nanotubes [13], boron nitride sheets [14] and fullerenes [15–17]. 2D materials, by virtue of their high surface area to volume ratio, offer greater benefits for H<sub>2</sub> storage. With the advent of nanotechnologies, unprecedented advances are being made in fabricating new 2D materials [18,19]. In recent years, several 2D materials have been demonstrated to be potential substrates for H<sub>2</sub> storage, e.g. metal decorated pristine graphene has shown great promise as a potential candidate for H<sub>2</sub> adsorption and efficient storage [20]. Recently, Li-decorated porous graphene and naphthalene were investigated as potential substrate for H<sub>2</sub> storage [21,22]. It was found that metal decoration enhances the capacity of H<sub>2</sub> storage of the substrate. Nevertheless, it is also well known that the adsorption of metallic atoms over pristine graphene is rather weak with binding energies to the tune of 0.6  $eV/H_2$  and as a result, metallic adatoms tend to cluster over the graphene substrate [23].

Phosphorene, a new p-type semiconducting 2D material of phosphorous, has been fabricated recently [24,25]. Phosphorene is unique compared to graphene because of its layerdependent bandgap which ranges between 0.3 eV (bulk) and 1.5 eV (monolayer), and therefore, has the capacity to perform in applications where gapless graphene and large band gap transition metal dichalcogenides (TMDSs) cannot cater to [24]. Phosphorene has a six-ringed structure, however, unlike graphene, it has a puckered structure due to sp3 hybridized P-P bonds resulting in a higher surface-area-to-volume ratio [26,27]. Moreover, although heavier than graphene, monolayer black phosphorous (MBP) is lighter than most of the known 2D materials [27] and can be exfoliated mechanically [28]. Recent research articles have demonstrated phosphorene to be promising in several applications, e.g. nanoelectronics, optoelectronics and photovoltaics [25,29-31]. Additionally, phosphorene has been recently investigated to evaluate its capability to adsorb metals [32], nonmetals [33], small molecules [34] and  $H_2$  [27,35]. Li et al. [27] have shown that pristine phosphorene is a rather weak substrate for the adsorption of  $H_2$  molecules with a binding energy of 0.07 eV/ $H_2$ , however, Lidecoration increased the binding energy to 0.18 eV. Yu et al. [35] reported that the phosphorene structures used by Li et al. [27] were unstable, and the gravimetric density in Li-decorated pristine phosphorene was limited to 4.4% [35].

We utilized defect engineering and Li-decoration to enhance the H<sub>2</sub> storage capacity of phosphorene. Li adatoms were selected owing to their light weight, which enhances the gravimetric density of the system and that it binds strongly with defective phosphorene substrates [36]. Previous researches also indicate that metallic adatoms tend to cluster on defect free substrates instead of homogeneously dispersing owing to their large cohesive energy [37,38]. This phenomenon severely limits the hydrogen storage performance of metaldecorated 2D materials. One way of tackling this problem is to introduce lattice defects in the substrate which would enhance the bond strength between the metallic adatom and the substrate, thereby inhibiting clustering of the adatoms [39,40]. This method is advantageous because lattice defects are inadvertently created during the synthesis of 2D materials. On the other hand, increased interaction between the adatom and the substrate reduces the amount of charge available to facilitate the interactions with the  $H_2$  molecules [41]. Therefore, we must find structural defects that would prevent clustering of adatoms without negotiating with large and efficient H<sub>2</sub> storage. Three defects, namely single vacancy (SV), double vacancy (DV) and Stone-Wales (SW) were considered as they are relatively stable and posses low formation energies [42]. In order to gain further insight into  $H_2$  adsorption onto defective phosphorene, differential charge density (DCD) and projected density of states (PDOS) were computed.

#### Material and computational detail

Plane wave pseudopotential based DFT calculations were performed using the open source package Quantum Expresso available under the GNU license [43]. For the DFT calculations, Perdew-Burke-Ernzerhof (PBE) exchange correlation functional [44] and Martins-Troullier NORMCONS type pseudopotential [45] in semilocal form generated by scalar relativistic calculation were used with Van der Waals correction in the dispersion corrected DFT (DFT-D) framework. In the DFT-D method, the van der Waals correction depends only on the positions of the nuclei and the exchange correlation functional and describes the nonlocal dispersion effects more accurately [46,47]. To avoid any inter-layer interactions between phosphorene surfaces from long-range van der Waals forces due to periodicity, we used a vacuum spacing of 20 Å. The cutoff energy for wave function expansion was set at 60 Ry (1 Ry = 13.61 eV) and that of the charge density was set at 480 Ry. Brillouin zone integrations were performed using a Monkhorst-Pack grid with  $4 \times 4 \times 1$  k-points [48]. The Methfessel–Paxton smearing with a degauss value of 0.02 was used in the simulations [49]. The binding energy of H<sub>2</sub> for each case can be calculated as:

$$E_{b} = \left(E_{phos+iH_{2}} - E_{phos} - E_{iH_{2}}\right)/i, \tag{1}$$

$$\mathbf{E}_{b} = \left(\mathbf{E}_{phos+Li+iH_{2}} - \mathbf{E}_{phos+Li} - \mathbf{E}_{iH_{2}}\right)/i, \tag{2}$$

where  $E_b$  represents the average binding energy of the hydrogen and  $E_{phos+iH2}$ ,  $E_{phos}$ ,  $E_{iH2}$  are the total energy of phosphorene and  $iH_2$  molecules, only phosphorene, and only  $iH_2$  molecules, respectively. Similarly,  $E_{phos+Li+iH2}$  and  $E_{phos+Li}$  represent the total energy of the full Li-decorated  $H_2$ -adsorbed phosphorene, and Li-decorated phosphorene, respectively. As per our definition, a negative binding energy implies that the adsorption is energetically favorable and occurring, while a positive binding energy indicates no adsorption. Smaller values of  $E_b$  imply a stronger binding of the  $H_2$  with the substrate.

In order to compute the charge density, self consistent field (SCF) calculations were performed using a convergence threshold of  $1e^{-5}$  Ry for the total system energy. The charge density differences ( $\Delta \rho$ ) have been computed using the relations:

$$\Delta \rho = \rho_{\text{phos}+\text{iH}_2} - \rho_{\text{phos}} - \rho_{\text{iH}_2},\tag{3}$$

$$\Delta \rho = \rho_{\text{phos}+\text{Li}+\text{iH}_2} - \rho_{\text{phos}+\text{Li}} - \rho_{\text{iH}_2}, \tag{4}$$

where  $\rho_{phos+iH2}$ ,  $\rho_{phos}$ , and  $\rho_{iH2}$  represent the charge density of phosphorene and  $iH_2$  molecules, only phosphorene, and  $iH_2$  molecules, respectively. For the Li decorated system,  $\rho_{phos+Li+iH2}$  and  $\rho_{phos+Li}$  represent the charge densities of Li-decorated phosphorene substrate with adsorbed  $H_2$  molecules and only Li-decorated phosphorene substrate, respectively.

At first, the defective phosphorene substrates were relaxed at their ground state. DFT calculations were performed with several initial positions of the  $H_2$  molecule, and the smallest

Table 1 — Formation energies for different defects in phosphorene and those reported in literature and comparison with those of graphene.								
Phospho	Graphene							
Defect	E <sub>f</sub> (eV)	Zhang et al. [36]	Hu and Yang [42]	[50,51]				
SV	1.71	1.29	1.62	7.38–7.85				
DV	1.46	1.28	1.90	7.52-8.70				
SW	1.11	1 63	1.01	4.50-5.30				

binding energies are reported here. In order to study the effect of Li-decoration, the Li atom was adsorbed over the defective phosphorene substrate. The Li-decoration was performed by choosing the favorable sites for Li, following the literature [36].  $H_2$  molecules were then added to the supercell and stabilized energetically.

## **Results and discussions**

As a benchmark, the binding energy associated with the adsorption of a  $H_2$  molecules on defective phosphorene substrates (with and without Li decoration) are compared with that of pristine phosphorene which has been reported to be -0.07 eV [27].

## Defective phosphorene

Similar to graphene, phosphorene is known to have stable structures with single vacancy (SV), double vacancy (DV) and Stone Wales (SW) defects as evidenced by their lower formation energies [42]. In this work, these three defects, i.e. SV, DV and SW were considered in order to study the H<sub>2</sub> storage capacity by defect engineering. Initially, when the pristine phosphorene supercell was created, the lattice parameters of pristine phosphorene were obtained to be 4.62 Å in the armchair direction and 3.30 Å in the Zigzag direction respectively. The defective phosphorene substrates were then relaxed at their ground state by energy minimization. Following the work by Zhang et al. [36], who have studied the stability and kinetics of a host of point defects in monolayer phosphorene, energetically stable single and double vacancy phosphorene substrates were created by removing P atoms from their lattice positions and Stone-Wales defect by rotating a P–P bond by 90° with regard to the midpoint of the bond.

The formation energies of the defective phosphorene have been calculated using the relation [36],

$$E_{f} = E_{defect} - \frac{N - n}{N} E_{pristine},$$
(5)

where  $E_{defect}$ ,  $E_{pristine}$  are the total energies of the defective and defect-free (pristine) phosphorene, respectively, N is the number of atoms in the pristine unit cell, and *n* is the number of atoms removed from the pristine cell to create the defect (*n*=1 for SV and SW, and *n*=2 for DV). As summarized in Table 1, the calculated formation energies for the DV 5-8-5, SV



Fig. 1 – Schematic models of Pristine, Single vacancy (SV), Double vacancy (DV) and Stone Wales (SW) defective phosphorene.



Fig. 2 – Binding energies ( $E_b$ ) of  $H_2$  molecules for pristine [27], Single Vacancy (SV), Double Vacancy (DV), and Stone Wales (SW) defective phosphorene and Li-decorated defective phosphorene substrates.

Table 2 – Binding energy ( $E_{h}^{H_2}$ ), perpendicular height (Z)							
from substrate, and bond length (R) of H <sub>2</sub> and binding							
energy of Li (E <sup>Li</sup> <sub>h</sub> ) in different defective and Li-decorated							
defective phosphorene. The binding energies correspond							
to adsorption of one H <sub>2</sub> molecule. The perpendicular							
height is measured from the closest phosphorous atom							
in defective phosphorene and from Li atom in Li-							
decorated defective phosphorene.							

System	$E_b^{H_2}$ (eV)	Z (A)	R (A)	$E_b^{L1}$ (eV)
SV	-0.058	2.02	0.745	-
DV	-0.062	2.95	0.746	-
SW	-0.031	2.27	0.744	-
SV-Li	-0.48	1.82	0.750	-3.21
DV-Li	-0.23	1.88	0.751	-3.02
SW-Li	-0.28	1.82	0.750	-2.49

5-9, and SW 55-77 defects are 1.46 eV, 1.71 eV, and 1.11 eV, respectively. These magnitudes of defect formation energies are in-sync with those reported in literature (Table 1). The formation energies for SV, DV and SW defective phosphorene are listed in Table 1 and are compared to those reported in

literature. For comparison, the formation energies of graphene from literature are also listed. It can be noted that formation of defects in phosphorene is easier than in graphene. The optimized structure of these point defects in phosphorene obtained from our DFT simulations are shown in Fig. 1.

#### H<sub>2</sub> adsorption on defective phosphorene

The binding energies obtained by using Eq. (1), are presented in Fig. 2. These values are also presented in Table 2 along with the distance between the H<sub>2</sub> molecule and the nearest phosphorous atom in the relaxed structure. It is observed that, among the defective phosphorene substrates, the DV phosphorene has the best binding energy  $(-0.062 \text{ eV/H}_2)$ , which is well above the acceptable limit for effective  $H_2$  storage. In order to showcase the interaction between the H<sub>2</sub> molecule and the phosphorene substrates, the differential charge density (DCD) and partial density of state (PDOS) are depicted in Figs. 3 and 4. It can be noted from the charge density difference plots that the charges are spread apart from the  $H_2$ molecules indicating interactions between the substrates and the H<sub>2</sub> molecules. The change in the PDOS of electrons in the p shell of the H<sub>2</sub>-adsorbed defective phosphorene as compared to that in the phosphorene substrate without H<sub>2</sub> shows the charge transfer from the H<sub>2</sub> molecules to the substrates (Fig. 4(a)). The PDOS of the s electron in  $H_2$  reduces after adsorption as compared to that of only  $H_2$  (Fig. 4(b)).

#### H<sub>2</sub> adsorption on Li-decorated defective phosphorene

DFT simulations of  $H_2$  adsorption onto Li decorated defective phosphorene were performed to evaluate the effects of Li decoration onto  $H_2$  adsorption. At first Li atoms were adsorbed over the defective phosphorene substrates by choosing favorable sites reported previously.  $H_2$  molecules were then added to the supercell and stabilized energetically. Similar to literature, in our simulations, enhanced interactions were also observed between Li adatoms and the defective phosphorene substrates. For example, the binding energies



Fig. 3 – Differential Charge Densities (DCDs) of in  $H_2$  adsorbed (a) SV 5-9, b (b) DV 5-8-5, and (c) SW 55-77 phosphorene. The P and H atoms are represented by violet and blue spheres. Yellow and light blue regions indicate charge accumulation and depletion, respectively. The isosurface level is set to be  $9e^{-5} e \text{ Å}^{-3}$  for (a) SV, (b) DV and (c) SW. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)



Fig. 4 – PDOS of (a) P and (b) H atoms before and after H<sub>2</sub> adsorption with Fermi energy shifted to zero. SV represents only SV phosphorene, H<sub>2</sub> represents only H<sub>2</sub>, and SV+H<sub>2</sub> represents SV phosphorene and H<sub>2</sub>.



Fig. 5 - Schematic models of Li-decorated SV, DV and SW defective phosphorene.



Fig. 6 – Differential Charge Densities (DCDs) in H<sub>2</sub> adsorbed Li-decorated defective phosphorene. The red sphere represents Li atom and dark blue spheres represent H<sub>2</sub> atoms. The isosurface levels are set to 0.00045 e Å<sup>-3</sup>. Yellow and light blue regions indicate charge accumulation and depletion, respectively. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)



Fig. 7 – PDOS of SV phosphorene (Fermi energy shifted to zero) with and without Li.

of Li adsorption onto SV, DV and SW phosphorene were calculated to be -3.21 eV, -3.02 eV, and -2.49 eV, respectively. These binding energies are much stronger than the cohesive energy of bulk Li. The binding energies of Li adsorption on defective phosphorene are in agreement with previous reports [52]. The atomic configurations of  $H_2$  adsorbed on Lidecorated defective phosphorene are shown in Fig. 5. Following the relaxation of Li-decorated phosphorene, a  $H_2$  molecule was added to the relaxed Li-decorated defective

phosphorene system, in such a way that one of the H atoms in the  $H_2$  molecule faced towards the Li atom in the vertical direction. Our DFT calculations predict that Li-decoration significantly enhances the  $H_2$  adsorption capacity of defective phosphorene. Li-decorated SV defective phosphorene was found to yield the best binding energy (Eq. (2)), as strong as  $-0.48 \text{ eV}/H_2$ , and the SW and DV defective phosphorene were predicted to have binding energies of  $-0.28 \text{ eV}/H_2$ , and  $-0.23 \text{ eV}/H_2$ , respectively (Fig. 2, Table 2). The vertical distance between the  $H_2$  molecules and the Li atoms are also listed in Table 2. It is found that in all the relaxed structures the distance between the  $H_2$  molecules and the Li atoms lie in the range between 1.82 Å and 1.88 Å.

The differential charge densities (DCD) were computed from the DFT results for understanding the charge transfer among different constituent atoms in the presence of *Li* adatoms (Fig. 6). The DCDs show regions of charge accumulation and depletion on either side of the hydrogen molecules indicating strong polarization of the hydrogen molecules. The charge density suggests that the hydrogen binds by a weak electrostatic dipole mechanism without any strong hybridization or covalent bonding [20]. From these results, we deduce that the H<sub>2</sub> molecules were adsorbed by the *Li* atom and there were no Kubas interactions [41]. The weaker bonding type is believed to be better in the context of reversible H<sub>2</sub> storage where less energy can activate the release of stored H<sub>2</sub>.



Fig. 8 – PDOS of Li in Li-decorated defective phosphorene (left) before and (right) after H<sub>2</sub> adsorption.



In order to investigate the effect of Li on the charge transfer mechanism leading to better H<sub>2</sub> adsorption, the projected density of states were compared for SV phosphorene and Li-decorated SV phosphorene. Fig. 7 depicts the total PDOS of s and p shells of electrons in phosphorene with and without Li decoration. It can be seen that the PDOS of p shell electrons in phosphorene is affected (around 3 eV) due to the presence of Li indicating interaction between them. The PDOSs of the valence electrons belonging to Li in Li-decorated defective phosphorene before and after H<sub>2</sub> adsorption are shown in Fig. 8, to understand the charge transfer from the Li atom. A comparison of the PDOSs of Li in defective phosphorene before and after H<sub>2</sub> adsorption shows that the 2s and 2p orbitals of Li in defective phosphorene have higher peaks as compared to those when H<sub>2</sub> is adsorbed in Li-decorated defective phosphorene, indicating enhanced charge transfer from Li to H<sub>2</sub>. The fact that the intensity of the peaks reduces the most for the SV defected phosphorene and least when interacting with DV defected phosphorene corroborates with the prediction of lowest binding energy in SV defective

phosphorene compared to DV defective phosphorene (Figs. 9 and 10).

## Gravimetric density in Li-decorated defective phosphorene

The gravimetric density was calculated in order to estimate the maximum capacity of  $H_2$  storage in Li-decorated SV phosphorene. We focused on SV phosphorene since the Lidecorated phosphorene with SV defect showed the suitable binding energy (Section  $H_2$  adsorption on Li-decorated defective phosphorene). The gravimetric density of  $H_2$  storage is given by

$$C_{g} = \frac{pm_{H_{2}}}{qm_{P} + rm_{Li} + pm_{H_{2}}},$$
 (6)

where  $m_{H2}$ ,  $m_P$ ,  $m_{Li}$  and p, q, r represent the atomic mass and number of  $H_2$ , P, and Li, respectively. In order to maximize  $H_2$ adsorption on the phosphorene substrate, it was decorated by two Li atoms on both sides of the substrate followed by  $H_2$ adsorption. To study the gravimetric density, the number of



Fig. 10 – Charge density of Li-decorated SV phosphorene. The violet, dark blue and red spheres represent P, H and Li atoms, respectively. The isosurface levels are set to 0.002 e  $Å^{-3}$  to show the charge density differential around the Li atom. Yellow and light blue regions indicate charge accumulation and depletion, respectively. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)



Fig. 11 – Binding energy (black) and gravimetric density (blue) of  $H_2$  adsorption in of Li-decorated SV phosphorene. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

 $H_2$  molecules was sequentially increased and the average binding energy has been computed. It can be noted that after adsorption of the first H<sub>2</sub> molecule, the consequent molecules will have a lesser sequential binding energy as reported by Yu et al. [35]. However, when considering gravimetric density, the average binding energy computed from Eq. (2) is considered following literature [27,35,53]. In both cases (Li atom on one side and both sides), the number of  $H_2$  molecules has been increased till the average binding energy decreased to around  $-0.1 \text{ eV/H}_2$ . It was also noted that the H<sub>2</sub> molecules do not move away from the substrate indicating adsorption to phosphorene. It can be noted that defect density was not varied in this study. The average binding energy and the corresponding gravimetric density of H2 adsorption over SV phosphorene decorated with one Li atom are depicted in Fig. 11(a) as a function of number of molecules of  $H_2$  adsorbed. Additionally, in Fig. 11(b), the average binding energy and gravimetric density associated with H<sub>2</sub> adsorption for Li decoration on both sides of SV phosphorene are presented. As can be observed from the figures, an increase in the number of  $H_2$  molecules affects the average binding energy, leading to a reduced capability for further adsorption. In case of SV phosphorene decorated with one Li atom (Fig. 11(a)), a gravimetric density of 3.6% was predicted with a binding energy being around  $-0.1 \text{ eV/H}_2$  (p=35, q=1, r=20 in Eq. (6)). Li decoration on both sides of the SV phosphorene substrate (Fig. 11(b)) significantly increased the gravimetric density to 5.3% with a binding energy of around -0.1 eV (p=35, q=2, r=30 in Eq. (6)). The gravimetric density of  $H_2$  in metal decorated defective graphene was reported to be 7.02% [54], higher than this, however, it must be noted that the formation energies of the defects in graphene is much higher than those in phosphorene (Table 1).

# Conclusions

Density Functional Theory (DFT) calculations were performed to investigate the viability of  $H_2$  storage in defective and Li-decorated defective phosphorene. Three types of defects, namely the SV 5-9, DV 5-8-5 and SW 55-77, were studied due to their stability and low formation energies. The DFT calculations suggest that the defects in phosphorene do not enhance the H<sub>2</sub> storage capacity compared to that of pristine phosphorene. The binding energies in SV, DV and SW defective phosphorene were computed to be around 0.03-0.06 eV/H<sub>2</sub>, well below the acceptable range for reversible H<sub>2</sub> storage. Further, the defective phosphorene substrates were decorated with Li atoms, and the capacity of  $H_2$ storage in Li-decorated defective phosphorene substrates was studied. It was observed that Li-decoration in defective phosphorene significantly enhances the H<sub>2</sub> storage, with a binding energy as low as -0.48 eV/H<sub>2</sub> for Li-decorated SV defective phosphorene. DCDs and PDOSs evidence the unique feature of the Li decoration in terms of the charge transfer from Li. The weak bonding of H<sub>2</sub> atoms with the Lidecorated phosphorene makes the H<sub>2</sub> storage and release favorable for effective and reversible H<sub>2</sub> adsorption. In order to study the maximum capacity of H<sub>2</sub> storage, the Lidecorated SV phosphorene (the best case among the three defects) was further investigated by adding more H<sub>2</sub> molecules to determine the binding energy and gravimetric density. To do that, the SV defective phosphorene was decorated with Li on both sides in order to increase the H<sub>2</sub> adsorption capability. A gravimetric density of 3.6% and 5.3% could be achieved by single sided and double sided Li decoration on SV phosphorene, respectively. This is higher than the 4.4% gravimetric density obtained for Li-decorated pristine phosphorene [35] indicating that Li-decoration in defective phosphorene enhances H<sub>2</sub> adsorption better than its pristine counterpart.

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